

# Quantum Computation: Introduction

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# From classical physics to quantum

- **Classical physics**: Developed over centuries (Archimedes, Newton, Maxwell)
- Objects have well-defined properties, independent of how they are measured
- **Quantum mechanics**: First half of 20th century (Planck, Einstein, Bohr, Schrödinger, Heisenberg)
- One of our best physical theories, never been contradicted by experiment
- Not just in the lab: 1/3 of our GDP depends on quantum
- Many “weird” effects:  
superposition, interference, entanglement

# Quantum computers

- Current computers (in theory and practice) are based on **classical** physics
- Feynman, Benioff ( $\pm 1982$ ):  
What about **quantum mechanical** computers?  
Can we use those weird effects for useful computation?
- Deutsch ('85): universal quantum Turing machine
- Peter Shor: efficient algorithm for factoring ('94)
- Since then: fast growing field
  1. **can we build it?**
  2. **what can it do?**
- We focus on second question: quantum algorithms

# Overview of the course

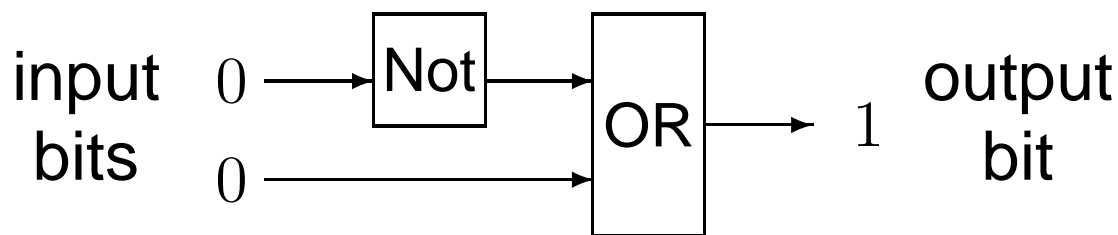
- Quantum computation: **introduction** (today)
- Quantum computation: **main algorithms**
- Quantum **communication**

# Overview of this lecture

- What are classical algorithms?
- What are quantum algorithms?
- Simple quantum algorithms:
  - Deutsch-Jozsa
  - Simon

# Classical algorithms

- Operate on **bits**
- Two main models: Turing machines, Boolean circuits
- **Circuits** are easier to generalize to quantum
- Directed acyclic graph of AND, OR, NOT **gates**
- Starting nodes:  $n$  input bits, additional workspace
- This computes some function by evaluating all gates



- **Efficient** computation: **polynomial-size** circuits

# From classical to quantum

- bits  $\longrightarrow$  qubits
- AND/OR/NOT gates  $\longrightarrow$  unitary quantum gates
- classical circuit  $\longrightarrow$  quantum circuit
- reading the output bit  $\longrightarrow$  measuring final state

# Recap of linear algebra 1

- Vector space  $V$  over field  $\mathbb{F}$ : set of objects such that
  1.  $v, w \in V \Rightarrow v + w \in V$  (closed under addition)
  2.  $v \in V, a \in \mathbb{F} \Rightarrow av \in V$  (closed u. scalar multiplication)

Think:  $V = \mathbb{C}^d$ ,  $v = (v_1, \dots, v_d)^T$ , basis  $\{e_1, \dots, e_d\}$

- Inner product:  $\langle v|w \rangle = \sum_{i=1}^d v_i^* w_i$

- Orthogonal:  $\langle v|w \rangle = 0$

- Norm:  $\|v\| = \sqrt{\langle v|v \rangle} = \sqrt{\sum_{i=1}^d |v_i|^2}$

- Unit vector: norm 1



# Recap of linear algebra 2

- **Linear transformation**  $A : V \rightarrow W$

1.  $u, v \in V \Rightarrow A(u + v) = A(u) + A(v)$

2.  $v \in V, a \in \mathbb{F} \Rightarrow A(av) = aA(v)$

Think:  $V = W = \mathbb{C}^d$ ,  $A$  is  $d \times d$  matrix

- $A$  is **Hermitian** if  $A = A^*$  (conjugate transpose)

- $A$  is **unitary** if  $A^{-1} = A^*$

Equivalent:

$A$  is norm-preserving,  
columns of  $A$  are orthonormal

# Recap of linear algebra 3

- Tensor product of matrices:

$$A \otimes B = \begin{pmatrix} A_{11}B & \cdots & A_{1d'}B \\ & \ddots & \\ A_{d1}B & \cdots & A_{dd'}B \end{pmatrix}$$

- Special case: vectors  $\begin{pmatrix} a_1 \\ \vdots \\ a_d \end{pmatrix} \otimes \begin{pmatrix} b_1 \\ \vdots \\ b_e \end{pmatrix} = \begin{pmatrix} a_1b_1 \\ a_1b_2 \\ \vdots \\ a_db_e \end{pmatrix}$

- Tensor product  $V \otimes W$  of spaces  $V$  and  $W$ :  
take basis  $\{v_1, \dots, v_d\}$  for  $V$ , basis  $\{w_1, \dots, w_e\}$  for  $W$ ,  
then  $V \otimes W = \text{span}\{v_i \otimes w_j : 1 \leq i \leq d, 1 \leq j \leq e\}$   
Note: dimension of  $V \otimes W$  is  $d \cdot e$

# Quantum bits

- **Classical bit:** value 0 or value 1
- Basis states of a **2-dimensional vector space:**

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

- **Qubit:** superposition  $\alpha_0|0\rangle + \alpha_1|1\rangle = \begin{pmatrix} \alpha_0 \\ \alpha_1 \end{pmatrix} \in \mathbb{C}^2$

- We require  $|\alpha_0|^2 + |\alpha_1|^2 = 1$

- Examples:  $\frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle$   
 $\frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}e^{i\pi/4}|1\rangle$   
 $\sin(\alpha)|0\rangle + \cos(\alpha)|1\rangle$

# More qubits

- Two qubits: tensor product space, with 4 basis vectors

$$|0\rangle \otimes |0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}, |0\rangle \otimes |1\rangle, |1\rangle \otimes |0\rangle, |1\rangle \otimes |1\rangle$$

- Abbreviate  $|a\rangle \otimes |b\rangle = |a\rangle|b\rangle = |a, b\rangle = |ab\rangle$

- 2-qubit state:  $|\phi\rangle = \sum_{x \in \{0,1\}^2} \alpha_x |x\rangle \in \mathbb{C}^4$

- Example: EPR-pair:  $\frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle$  (entangled)

- $n$ -qubit state:  $|\phi\rangle = \sum_{x \in \{0,1\}^n} \alpha_x |x\rangle \in \mathbb{C}^{2^n}$

# Quantum states and dynamics

- $n$ -qubit state  $|\phi\rangle = \sum_{x \in \{0,1\}^n} \alpha_x |x\rangle = \begin{pmatrix} \alpha_{0\dots 0} \\ \vdots \\ \alpha_{1\dots 1} \end{pmatrix} \in \mathbb{C}^{2^n}$
- *Informally*: we are in all  $2^n$  basis states simultaneously
- *Formally*:  $|\phi\rangle$  is a **unit** vector in  $2^n$ -dimensional space
- Two kinds of quantum operations on  $|\phi\rangle$ :
  1. **Unitary transform** of the amplitude-vector
  2. **Measurement**

# Measurement

- Measuring quantum state  $|\phi\rangle = \sum_{x \in \{0,1\}^n} \alpha_x |x\rangle$   
gives  $|x\rangle$  with probability  $|\alpha_x|^2$ ; state **collapses** to  $|x\rangle$
- Note: probabilities sum to 1 because  $|\phi\rangle$  is a unit vector
- We can also measure **part** of a state.  
The state then collapses to the part that is “consistent” with the measurement outcome
- Example: measure 2nd register of  $\frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle |f(x)\rangle$   
gives  $|a\rangle$  with probability  $\frac{|\{x: f(x)=a\}|}{2^n}$ ;  
State collapses to  $\frac{1}{\sqrt{|\{x : f(x) = a\}|}} \sum_{x: f(x)=a} |x\rangle |a\rangle$

# Quantum gate: unitary on 1 or 2 qubits

• 1-qubit NOT gate:  $X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

• 1-qubit Hadamard gate:  $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$

• 1-qubit  $\pi/4$ -gate:  $T = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{pmatrix}$

• 2-qubit controlled-NOT:  $C = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$

# Example: fun with Hadamard

- $H|0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) = |+\rangle$

Measurement gives  $|0\rangle$  or  $|1\rangle$  with probability 1/2

- $H|1\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) = |-\rangle$

Measurement gives  $|0\rangle$  or  $|1\rangle$  with probability 1/2

- We can get **interference**:

$$H|+\rangle = \frac{1}{\sqrt{2}}(H|0\rangle + H|1\rangle) = \frac{1}{2}(|0\rangle + |1\rangle + |0\rangle - |1\rangle) = |0\rangle$$

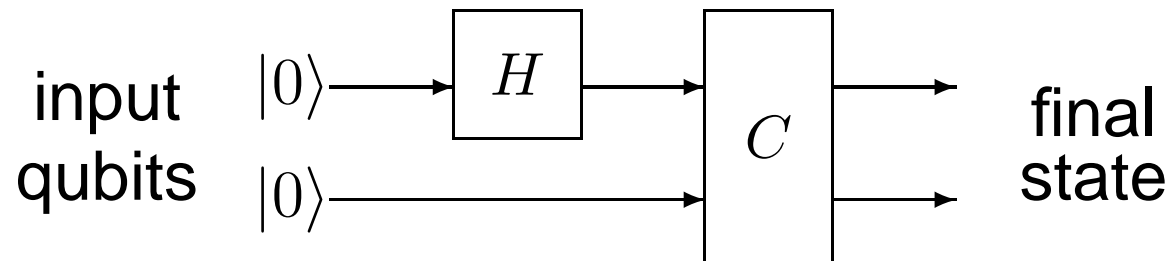
$$H|-\rangle = |1\rangle$$

- Hadamard on  $n$  qubits:  $H^{\otimes n}|x\rangle = \frac{1}{\sqrt{2^n}} \sum_{y \in \{0,1\}^n} (-1)^{x \cdot y} |y\rangle$



# Quantum circuits

- Circuit of gates transforms input state to **final state**



- Viewed as a big unitary:  $C(H \otimes I)$
- Final state:  $\frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$ , an EPR-pair
- **Measure** specific qubit of final state to obtain output
- $H$ ,  $T$ ,  $C$  gates can approximate any  $n$ -qubit unitary
- **Efficient** quantum computation: **polynomial-size** circuit

# Quantum parallelism

- Suppose classical algorithm computes  $f : \{0, 1\}^n \rightarrow \{0, 1\}^m$
- Then quantum circuit  $U : |x\rangle|0\rangle \mapsto |x\rangle|f(x)\rangle$  can compute  $f$  on all inputs **simultaneously!**

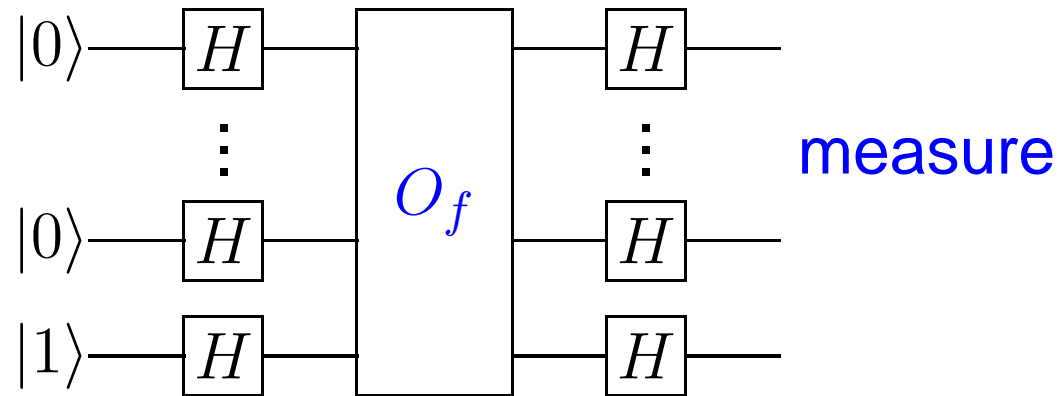
$$U \left( \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle|0\rangle \right) = \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle|f(x)\rangle$$

- This contains all  $2^n$  function values!
- But observing gives only one random  $|x\rangle|f(x)\rangle$
- All other information will be lost
- **More tricks needed for successful quantum computation**

# Deutsch-Jozsa problem

- Given: function  $f : \{0, 1\}^n \rightarrow \{0, 1\}$  ( $2^n$  bits), s.t.
  - (1)  $f(x) = 0$  for all  $x$  (**constant**), or
  - (2)  $f(x) = 0$  for  $\frac{1}{2} \cdot 2^n$  of the  $x$ 's (**balanced**)
- Question: is  $f$  constant or balanced?
- Classically**: need at least  $\frac{1}{2} \cdot 2^n + 1$  steps ("queries" to  $f$ )
- Quantumly**:  $O(n)$  gates suffice, and only 1 query
- Query: application of unitary  $O_f : |x, 0\rangle \mapsto |x, f(x)\rangle$
- More generally:  $O_f : |x, b\rangle \mapsto |x, b \oplus f(x)\rangle$  ( $b \in \{0, 1\}$ )
- Note:  $O_f |x\rangle |-\rangle = (-1)^{f(x)} |x\rangle |-\rangle$

# Deutsch-Jozsa algorithm



- Starting state:  $|\underbrace{0 \dots 0}_n\rangle|1\rangle$

- After first Hadamards:  $\frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle|-\rangle$

- Make one query:  $\frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} (-1)^{f(x)} |x\rangle|-\rangle$

- Forget about the  $|-\rangle$

# Deutsch-Jozsa (continued)

- After second Hadamard:

$$\frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} (-1)^{f(x)} \frac{1}{\sqrt{2^n}} \sum_{y \in \{0,1\}^n} (-1)^{x \cdot y} |y\rangle$$

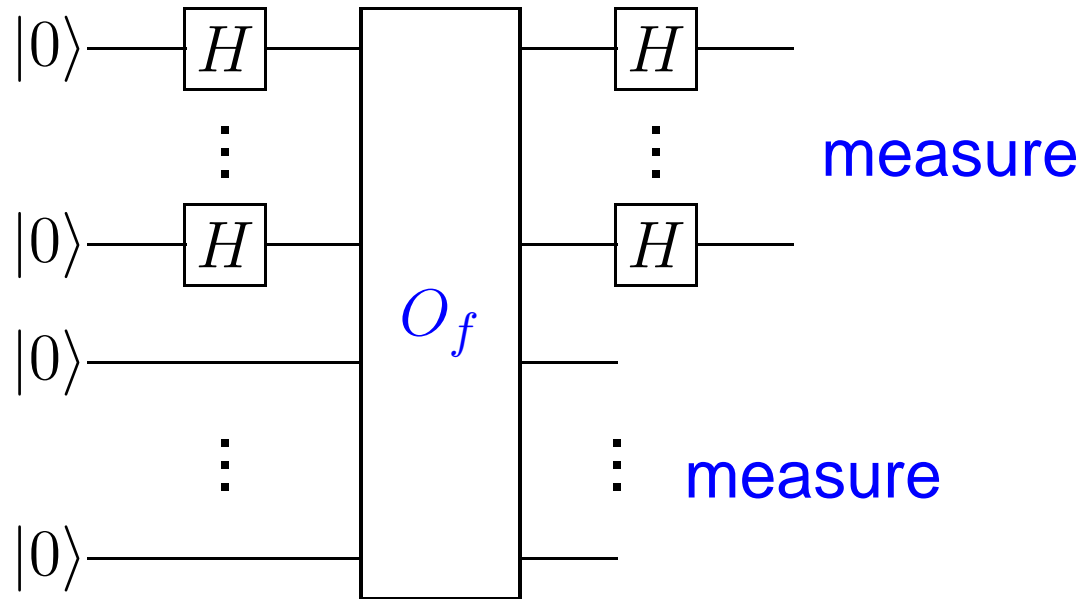
- $\alpha_{0\dots 0} = \frac{1}{2^n} \sum_{x \in \{0,1\}^n} (-1)^{f(x)} = \begin{cases} 1 & \text{if constant} \\ 0 & \text{if balanced} \end{cases}$

- Measurement gives right answer with certainty
- Big quantum-classical separation...
- But the problem is efficiently solvable by bounded-error classical algorithm (query  $f$  at a few random  $x$ )

# Simon's problem

- Given: function  $f : \{0, 1\}^n \rightarrow \{0, 1\}^n$  such that there exists  $s \in \{0, 1\}^n$  satisfying  $f(x) = f(y)$  iff  $x = y \oplus s$
- Note: if  $s = 0^n$  then  $f$  is a permutation (1-1), otherwise  $f$  is a 2-1 function
- Question: is  $s = 0^n$  or not
- Classically: need  $\sqrt{2^n}$  queries for high success prob
- Quantumly: solve in  $O(n)$  queries and  $O(n^3)$  gates
- Quantum algorithm is exponentially better, even compared with classical bounded-error algorithms

# Simon: quantum algorithm



- After  $H$ 's and  $O_f$ :  $\frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle |f(x)\rangle$
- Measure specific  $f(x)$ : 1st register  $\frac{1}{\sqrt{2}}(|x\rangle + |x \oplus s\rangle)$
- After  $H$ 's:  $\frac{1}{\sqrt{2^{n+1}}} \left( \sum_{y \in \{0,1\}^n} (-1)^{x \cdot y} |y\rangle + (-1)^{(x \oplus s) \cdot y} |y\rangle \right)$

# Simon's algorithm (continued)

- First  $n$  qubits:  $\frac{1}{\sqrt{2^{n+1}}} \sum_{y \in \{0,1\}^n} (-1)^{x \cdot y} (1 + (-1)^{s \cdot y}) |y\rangle$
- Note:  $|y\rangle$  has non-zero amplitude iff  $s \cdot y = 0 \bmod 2$
- Measure: get string  $y \in \{0,1\}^n$  s.t.  $s \cdot y = 0 \bmod 2$
- Repeat this  $2n$  times, giving  $y_1, \dots, y_{2n}$ , each with  $s \cdot y_i = 0 \bmod 2$
- W.h.p. there are  $n$  linearly independent  $y$ 's
- This system of linear equations  $s \cdot y_i = 0 \bmod 2$  determines  $s$  (solve via Gaussian elimination)
- Quantum algorithm uses  $O(n)$  queries and  $O(n^3)$  gates



# Classical lower bound

- Intuition: a classical algorithm can only query  $f$  at random points
- As long as it doesn't find a collision ( $x, y$  s.t.  $f(x) = f(y)$ ) it cannot distinguish 1-1 from 2-1 functions
- For uniform 2-1 function and fixed  $x, y$ :  
 $\Pr[f(x) = f(y)] \approx 1/2^n$
- With  $T$  queries, we have queried  $\binom{T}{2}$  specific pairs

$$\Pr[\text{see a collision}] \leq \text{Exp}[\# \text{collisions}] \approx \binom{T}{2} \frac{1}{2^n} \approx \frac{T^2}{2^{n+1}}$$

- If  $T \ll \sqrt{2^n}$  then algorithm can't distinguish 1-1 from 2-1
- Classical algorithm needs  $\approx \sqrt{2^n}$  queries

# Summary

- We introduced quantum mechanics
- We showed how to use it for computation:  
qubits, unitary gates, circuits, measurements
- Quantum algorithms can be better than classical  
(Deutsch-Jozsa and Simon)
- Next two lectures:
  - Main quantum algorithms: Shor and Grover
  - Quantum communication